Physics of complex systems and criticality

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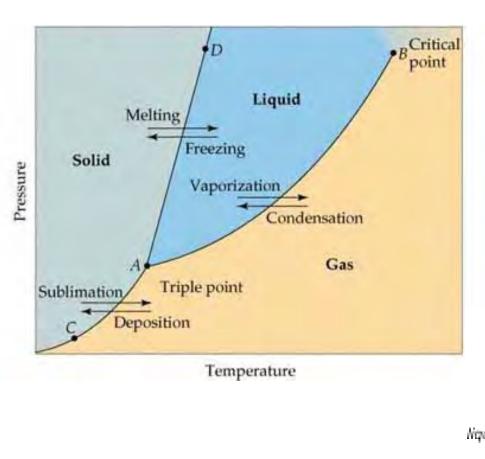


Diagram PT

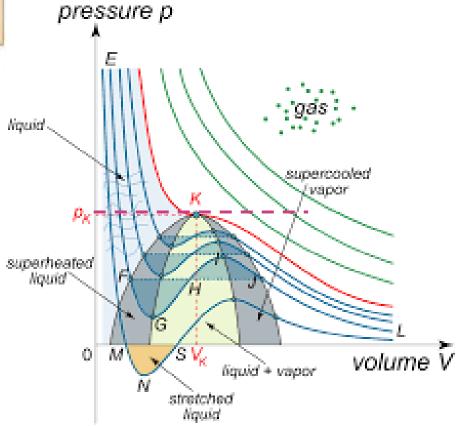
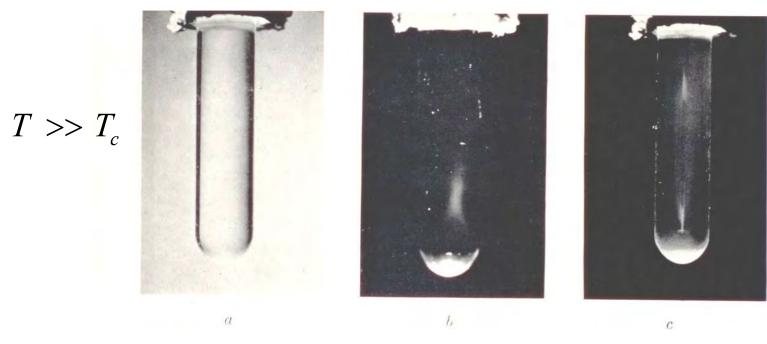
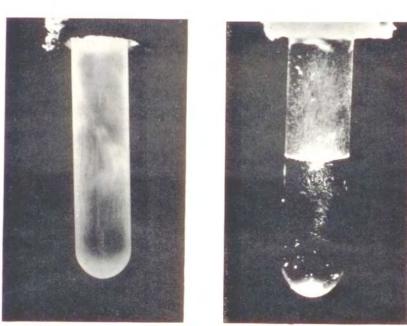


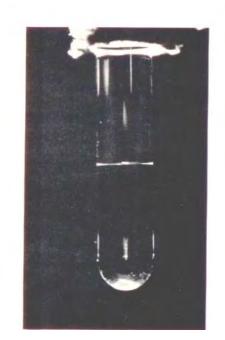
Diagram PV



critical opalescence

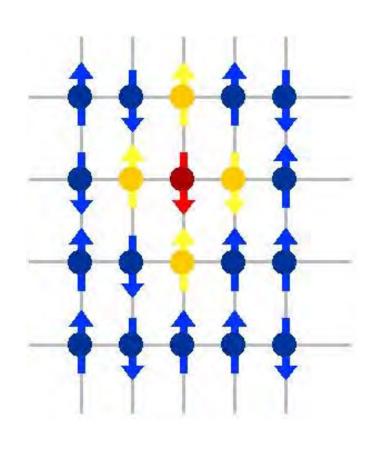
 $T \approx T_c$





 $T \ll T_c$

What is a critical state?



- Ising model: two-state spins
- Ferromagnetic interaction: the energy is lower if spins are parallel
- System evolves toward equilibrium state which minimizes the energy
- New state of a spin is assigned according to the probability

$$\propto e^{-\frac{E}{K_BT}}$$

- lacktriangle At $T\cong\infty$ all spins are uncorrelated
- ullet At T=0 all spins tend to allign

Emergence of spontaneous magnetization at T_c (h=0)

Symmetry breaking

$$H_{Ising} = -J \sum_{\langle ij \rangle} S_i S_j - h \sum_i S_i$$

For h=0 Hamiltonian of the Ising model is symmetric for spin inversion Configurations with magnetization +m and -m have the same probability

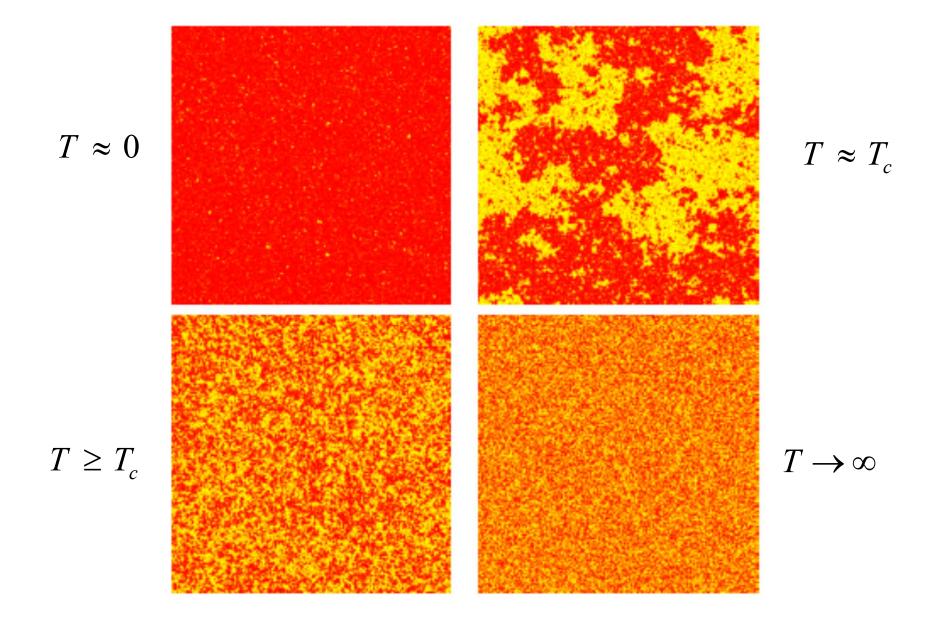


In the limit $N \to \infty$ the symmetry between the two phases can be broken and a spontaneous magnetization emerges

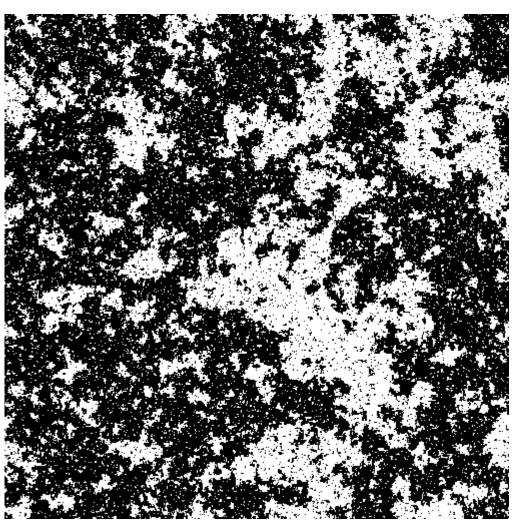
The system is trapped in a phase space subregion since the energy barrier is too high

Probability to escape

$$p \propto e^{-N^{(d-1)/d}}$$



- oclusters of all possible sizes are present
- •divergence of correlation range
- •divergence of fluctuations
- Self-similarity —— the largest cluster is fractal

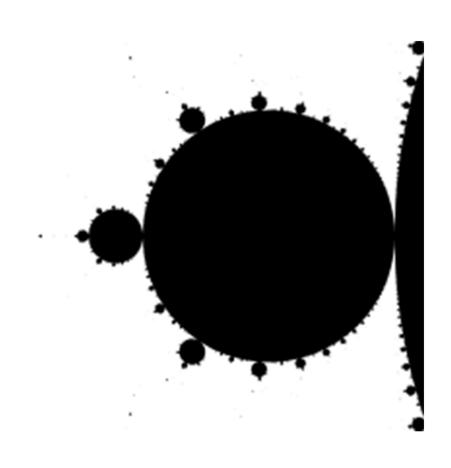


At the critical point
physical properties behave as
power laws:

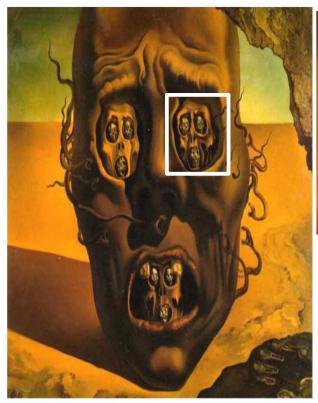
Order parameter
Response function
Specific heat

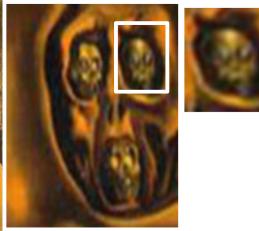
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Scale invariance - Mandelbrot set



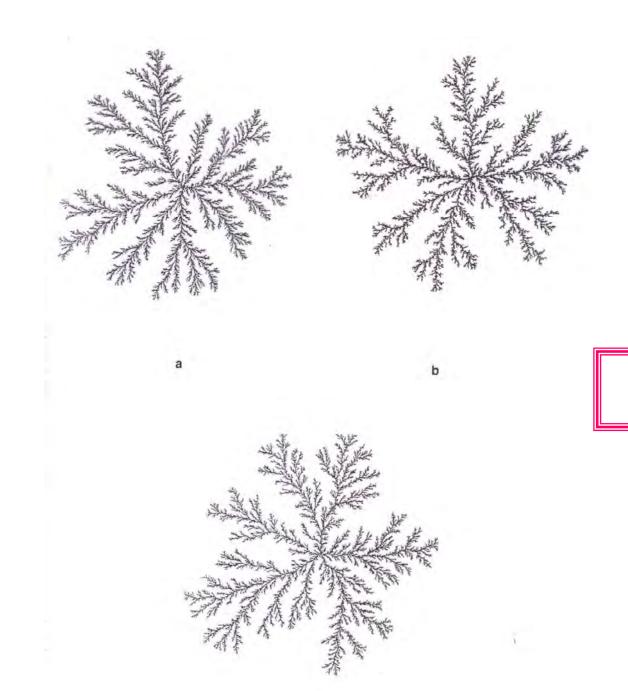
SELF-SIMILARITY





Salvador Dali "The war's face"

A part looks like the whole....



Self-similarity

Diffusion Limited Aggregation

Correlation function

The correlation function

$$G(r) = \langle m(r)m(0) \rangle - \langle m(0) \rangle^2$$

measures the correlations between fluctuations in 0 and r

The spatial extent of the correlations is called correlation length

The spatial integral of the correlation function provides the response function (derived by the fluctuation dissipation theorem)

$$\chi = \frac{1}{k_B T} \int d^d r G(r)$$

At the critical point

$$\xi \to \infty$$

$$\xi \to \infty$$

$$G(r) = G_0 \frac{k_B T}{\gamma} r^{-(d-2+\eta)} \exp(-r/\xi)$$

Experimentally can be measured by scattering experiments

Structure factor

$$S(q=0) \propto \tilde{G}(0) = \int G(\vec{r}) d\vec{r}$$

Critical Exponents

Reduced Temperature, t

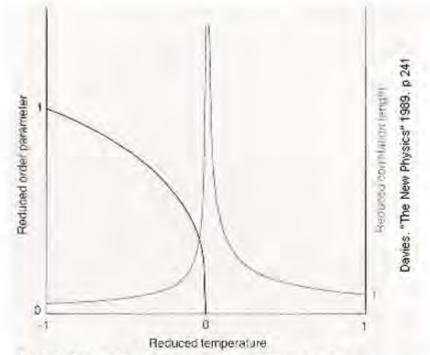
$$t \equiv \frac{T - T_C}{T_C}$$

Specific heat $C \propto |t|^{-\alpha}$

Magnetization $M \propto |t|^{\beta}$

Magnetic susceptibility $\chi \propto |t|^{-\gamma}$

Correlation length $\stackrel{z}{\xi} \propto \left|t\right|^{-\nu}$



Critical behavior of the order parameter an the correlation length. The order parameter vanishes with the power β of the reduced temperature t as the critical point is approached along the line of phase coexistance. The correlation length diverges with the power v of the reduced temperature.

The exponents display critical point universality (don't depend on details of the model).

Universality classes!

Universality 1.00 0.95 0.00 0.83 0.80 0:75 0.65 0.60 0.55 2/20

Densities of coexisting liquid and gas phases of a variety of substances, plotted against temperature, with both densities and temperatures scaled to their value at the critical point (From: E A Guggenheim, J. Chem. Phys., Vol. 13, 253, 1945)

Exponents are not all independent....

$$\alpha + 2\beta + \gamma = 2,$$

$$\gamma = \beta(\delta - 1),$$

$$2 - \alpha = \nu d,$$

$$\frac{\gamma}{\nu} = 2 - \eta.$$

A number of relations have been derived among them

PROBLEM: thermodynamics predicts them as INEQUALITIES

Whereas experimental values satisfy EQUALITIES!!!

Scaling - Widom hypothesis

- Near the critical point the main physical properties exibit power law behaviour
- Nice properties of power laws <u>invariant under rescaling!</u>

Suppose
$$y = f(x) = x^{\alpha}$$

Make the scale transformation

$$x \to x' = bx$$
 $y \to y' = cy$

Under rescaling

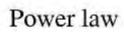
$$f(x) \rightarrow f(x') = cf(x'/b)$$

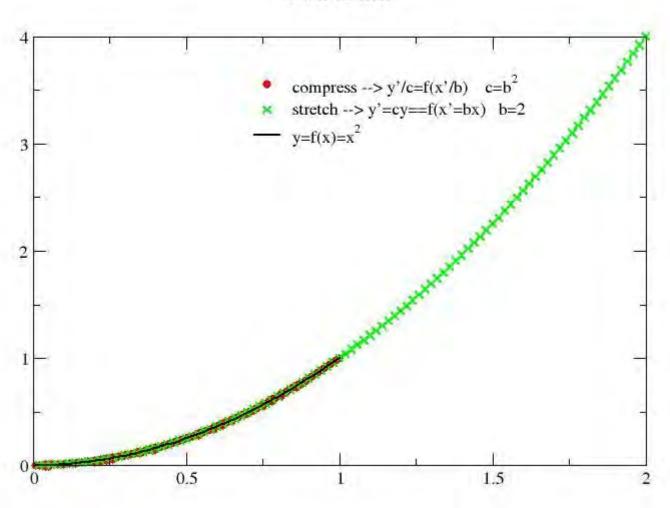
If
$$c = b^{\alpha}$$
 the function is invariant

Scale invariance

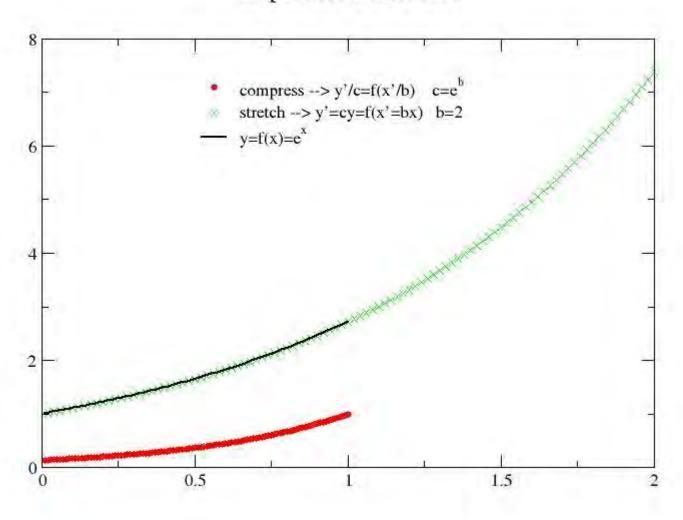
$$f(\lambda x) = g(\lambda)f(x)$$

homogeneous functions





Exponential function



For functions of more than one variable

$$f(\lambda^a x, \lambda^b y) = \lambda f(x, y) \quad \forall \lambda$$

Choose
$$\lambda = y^{-1/b}$$

Obtain

$$f(x,y) = y^{1/b} f(y^{-a/b}x,1) = y^{1/b} g(y^{-a/b}x)$$

Where g is an universal function

Assuming that thermodynamic potentials are GHF leads to relations among critical exponents as equalities

What is a complex system?

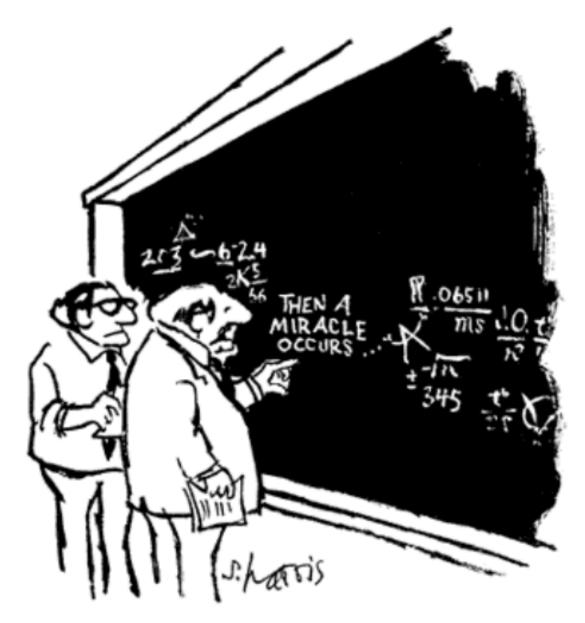
- Many components or degrees of freedom
- Interactions among components cooperative effects
- Emergence of «impredictable» macroscopic behaviour



SELF-ORGANIZATION

What is the signature of a complex system?

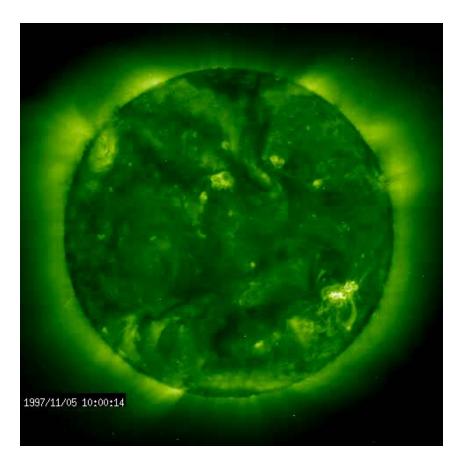
- Fundamental properties exhibit singular behaviour
- Emergence of power laws
- Absence of a characteristic scale

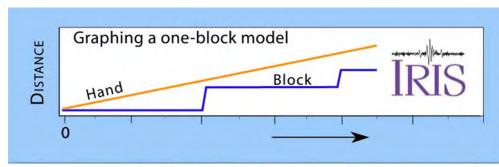


"I think you should be more explicit here in step two."

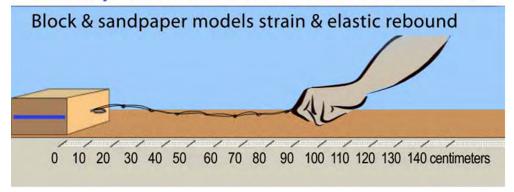
9. 8. 7. 1973. 1983. 1988. 1993. 1998. 2003.

Complex systems





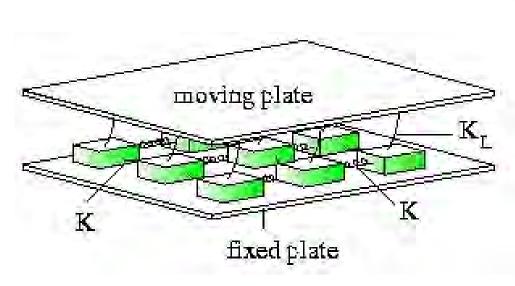
Earthquake Machine—Elastic Rebound

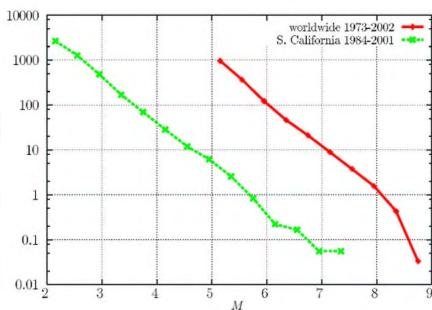


A single block pulled on a rough surface slips always the same distance

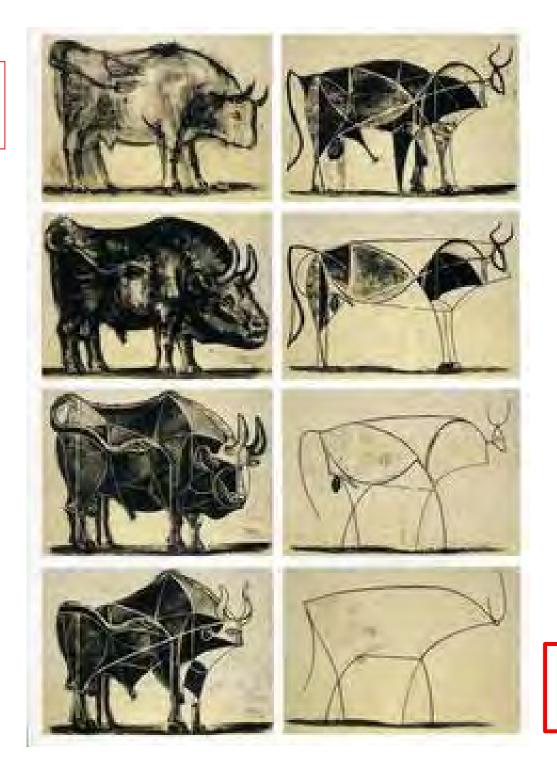
$$\Delta x = (\mu_s - \mu_d)(N/K)$$

On real faults earthquakes of all size are measured





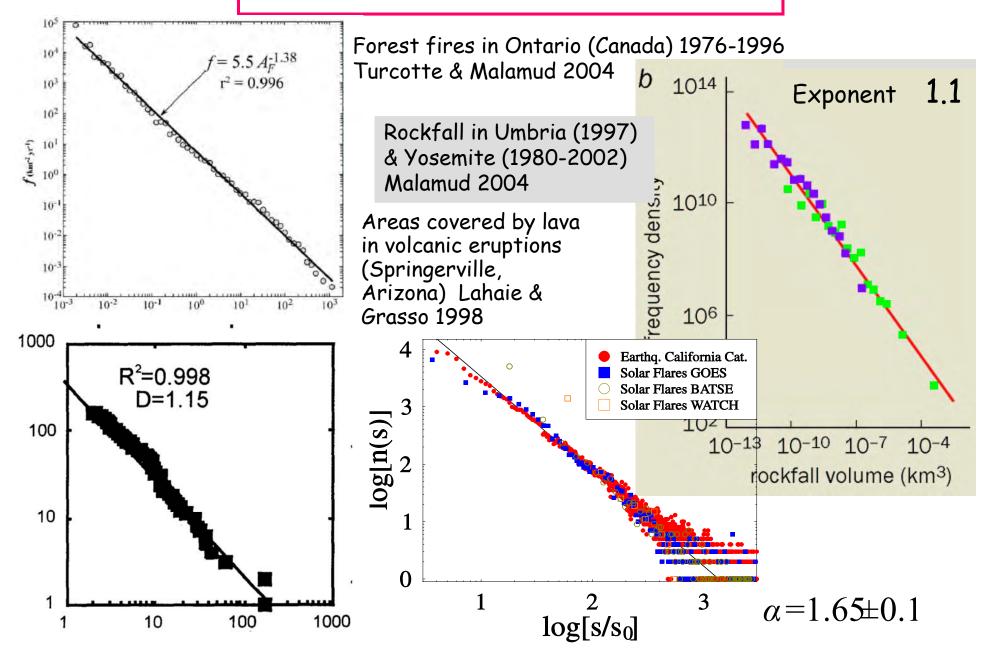
From complexity



Bull by Picasso

To Universality

Power laws in nature



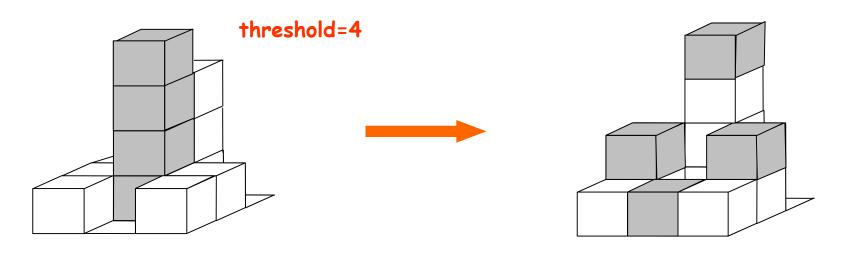
SELF-ORGANIZED CRITICALITY

Bak, Tang, Wiesenfeld, PRL 1987

Dynamical systems spontaneously evolving toward a critical state without parameter tuning _____ no characteristic event size

Sand pile

by adding at random one grain...



Size and duration distribution

$$P(s) \sim s^{-1}$$

$$P(s) \sim s^{-1}$$

 $P(T) \sim T^{-0.5}$

Fundamental ingredient: separation of time scales

- Slow scale: adding a grain
- Fast scale: propagation of an avalanche

SOC applied to many natural phenomena

- Slides and avalanches
- Neural activity
- Solar flares
- Fluctuations in confined plasma
- Biological evolution
- Earthquakes

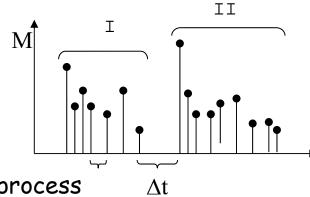
- The process generated by the sandpile or other standard SOC models is
 Poissonian absence of temporal correlations
- Additional ingredients must be introduced to generate a correlated process
- •In many stochastic processes in nature temporal correlations are present
- •How can we detect them?

Intertime distribution

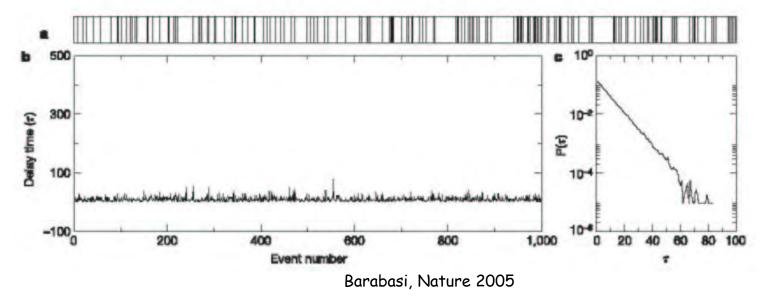
Probability distribution of intertimes

 Δi

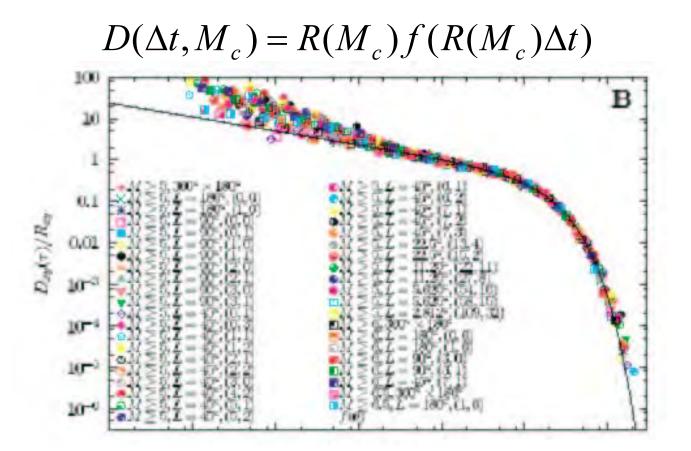
between consecutive events



ullet $P(\Delta t)$ is an exponential for a Poisson process



 It exhibits a more complex structure as temporal correlations are present in the process ightharpoonup Corral (PRL, 2004) rescaling Δt by the average rate in the area obtained a universal scaling law for the probability density



holds also for Japan, Spain, New Zeland... scaling function not universal (different areas are characterized by different rates)

Wiener - Khintchine Theorem

Autocorrelation function

$$K(s) = K(t_2 - t_1) = \langle A(t_2)A(t_1) \rangle$$

Power spectrum

$$S(f) = 4 \int_{0}^{\infty} ds K(s) \cos(2\pi f s)$$

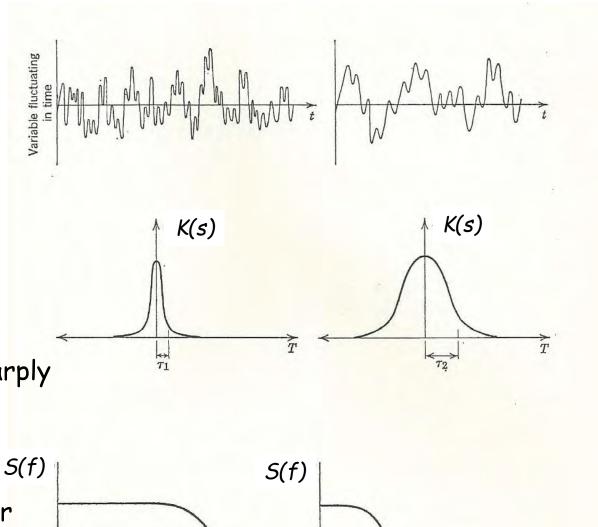
If the variable is very irregular (unpredictable) $K(s)=c\delta(s)$ Then and S(f)=2cfor all f

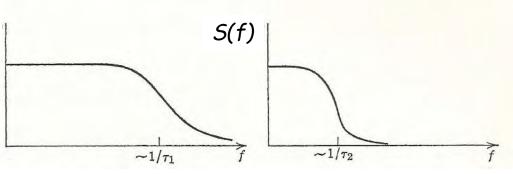
White noise

but K(0) would diverge!

In reality K(s) decays sharply within τ

Then S(f) is constant over a frequency range $1/\tau$





Color of noise

When not white noise is colored

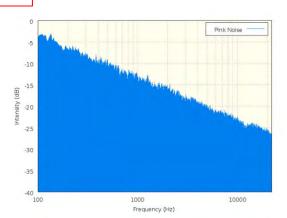
power law behaviour

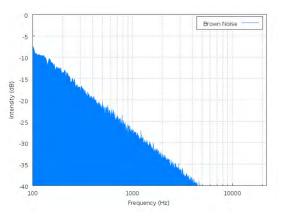
$$\int_{0}^{-\beta}$$

Pink (flicker) $\beta=1$

Brown (red) noise β =2

(by integrating white noise)





$$f^{-1}$$
 noise \longrightarrow long range temporal correlations

$$f^{-2}$$
 noise — uncorrelated signal

Take-home message

- Criticality implies the absence of a characteristic scale
- Emergence of power laws can be explained by SOC
- Power law distributions are not sufficient for criticality
- Necessity to verify the existence of long-range temporal correlations